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SEMI-ANNUAL PROGRESS REPORT

ON

ACTIVE CONTROL OF ROBOT MANIPULATOR COMPLIANCE

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SUMMARY

This progress report summarizes the research work performed at The Catholic University Of America on the research grant entitled "Active Control of Robot Manipulator Compliance, " (Grant#: NAG 5-780), supported by NASA/Goddard space Flight Center during the period of May 15th, 1986 to November 15th, 1986.

In this report we first present the modelling of the two-degree-of-freedom robot. Then the complete system including the robot and the hybrid controller is simulated on an IBM-XT Personal Computer. Simulation results showed that proper adjustments of controller gains enable the robot to perform successful operations. Further research should focus on developing a guideline for the controller gain design to achieve system stability.

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TABLE OF SYMBOLS

Values are given in SI Units.

- A vector from origin of compliant platform to
 i
 lower attachment points of actuator i with
 respect to the fixed coordinate system
- A known vector in moving coordinate system from i,m origin of moving coordinate system to lower attachment point of actuator i
- E known vector in fixed coordinate system from i crigin of fixed coordinate system to upper attachment point of actuator i
- d known distance between the upper attachment points of actuators 1 and 2
- F force
- f() vector function
- F vector of force error in Cartesian variables

F vector of desired force

f_ actuating forces

f. reaction force

g acceleration of gravity

1 inertia moment

[I] Identity matrix

J Jacobian

 $K_{\!\!\! E}$ stiffness of the reaction surface

 K_{rp} , K_{rf} proportional controller gains for position and force, respectively

κ_{rpd} derivative controller gain

l vector in fixed coordinate system from upper attachment point to lower attachment point of actuator i

L	Lagrangian
m,M	mass
P	potential energy
q i	generalized coordinates
C. i	generalized forces
q e	vector of position error in joint variables
Fi	vector in fixed coordinate system from origin of fixed coordinate system to origin of moving coordinate system
ř	vector in fixed coordinate system from origin of fixed coordinate system to lower attachment point of actuator i
S	compliance selection vector
Т	kinetic energy
стз	Euler angle transformation matrix

T _θ	torque
^{L(} rp	error in length due to the error in position
^U rf	error in length due to the error in force
V	velocity
Х _đ	vector of desired Cartesian position
Х _е	vector of position error in Cartesian variables
· X _e	vector of velocity error in Cartesian variables
y E	position of the reaction surface
α	unknown parameter vector which is vector of root of $f(\cdot)$
W	angular velocity
ů.	angular acceleration
ψ, θ, Φ	angular orientation of moving platform
Λ	position compensation function

1. INTRODUCTION

Robotics has many potential space applications that can be found in the space station research program [1], such as autonomous repair and replacement capability, automated maintenance, etc. One of the most essential robotic tasks is the assembly of large strucutures space. Product assembly in space such as mating and fastening of parts requires a very high precision. There however a tolerance problem that exists in assembly. The positional tolerance of many assembly tasks are quite high (thousands of inch), while those of manipulators are only of the orders of inch. To solve this problem, compliance can be provided to manipulator so that it can comply with the task-imposed constraint [2]. Compliant motion control is concerned with the control of a robot in contact with its environment Compliance can be provided passively by the deformation of a mechanical structure, such as mechanical spring or actively by a control system employing force sensors [4-5].

In many applications, there is a need for simultaneous control of position and force. There are two approaches to compliant motion control. Impedance or stiffness control

specifies dynamic relationships between force and position but do not specify desired forces or positions [3]. Hybrid control that controls position/orientation along specified degrees of freedom and independently controls force/torque along the remaining degrees of freedom [6]. The hybrid control is more suitable for planning a manipulator task and in this report we will present the development of a hybrid controller [7] for a six-degree-of-freedom robot recently built at NASA/GSFC [8] to study robotic assembly in space.

Since the dynamic equations of a six-degree-of-freedom robot is highly complicated, we start our design with a two-degree-of-freedom model. The results of simulated system will be presented and discussed.

2. THE NASA ROBOT

Recently a six-degree-of-freedom robot was built at the GODDARD SPACE FLIGHT CENTER (GSFC) to study the automated assembly of NASA hardware [8] for potential applications in the space station [1]. The robot consists mainly of a lower fixed platform, an upper movable platform and an intelligent end effector (IEE) (Fig. 1A). The movable platform is supported above the fixed platform by six axial rods that are extensible by recirculating ball screws. Six stepper motors were used to derive the ball screws and consequently the movable platform to provide

The IEE is attached the gross motion of the IEE. to another platform suspended from the movable platform by six spring-loaded pistons that provide passive compliance to the IEE by permitting strain on two opposing springs acting in the piston (Fig. 1B). Using force feedback, the robot has been able to perform several assembly tasks such as inserting a peg into the hole, screwing a bolt into a threaded hole. However, since the compliance was provided passively, the assembly process was very slow, for example, the robot took 3 minutes to complete the inserting of a peq into a hole. Therefore, research effort has been made to study the implementation of the hybrid control scheme into the robot system so that the IEE compliance can be provided actively by six electromechanical actuators replacing the existing spring-loaded pistons [10].

3. THE HYBRID CONTROL SCHEME

In many applications, it is necessary to control simultaneously position and force. Such control called hybrid control, that controls position /orientation along specified degrees of freedom and independently controls force/torque along the remaining degrees of The degrees of freedom are in a Cartesian freedom. coordinate system which is called the constraint frame and is denoted by the symbol [C] [11-13]. In a hybrid actuator control each signal controller scheme, composed of several components, one for each force

controlled degree of freedom in [C] and one for each position controlled degree of freedom. The actuator signal is computed by [7]:

$$T_{i} = \sum_{j=1}^{N} \{ \Gamma_{ij} [s_{j} \Delta^{f}] + \psi_{ij} [(1-s_{j}) \Delta^{X}] \}$$
 (1)

Where

 $T_i = torque$ applied by the ith actuator

 Δf_{ij} =force error in jth degree of freedom of [C].

 ΔX_{ij} =position error in jth degree of freedom of [C].

 Γ and ψ = force and position compensation functions, ij respectively, for the jth input and the jth output.

S =component of comliance selection vector.

The compliance selection vector S, is a binary N-tuple that specifies which degrees in [C] are under force control (sjj=1) and which are under position control (sjj=0).

Fig. 2 illustrates the implementation of the hybrid control scheme into existing NASA robot. The complete robot system consists of a position sensor, a six-degree-of-freedom force sensor, 12 controllers with 6 controllers for the position control loop and 6 controllers for the force/torque control loop, 6 linear actuators and the NASA robot. The compliance selection matrix is a (6x6) diagonal matrix.

Two feedback loops compute errors in both applied force/torque and actual position, with the lower loop

controlling the force/torge. The two independent feedback loops with 6 controllers, provide the drive signals to the actuators[10].

Current research does not prescribe particular control laws for the regulation of errors [7]. Therefore, our research objective is to make a comparative evaluation of different types of control laws in order to select an optional set of controllers for a particular robotic assembly task. Doing this a guideline can be developed for the design of these controllers. The controllers and the coordinate transformations are then implemented on an appropriate computer.

4. COORDINATE TRANSFORMATIONS

There are two coordinate transformations , one for transforming coordinate from Cartesian space into joint space, and one from joint to Cartesian space. With the arrangement of the linear actuators as seen on Fig. 1, the robot does not have independent drive systems for each degree of freedom. The precision motion is achieved in all degrees of freedom by a combination of actuator extension. Therefore, a transformation transforming desired position and orientation of the IEE into actuator extensions is needed. This problem is well-known as an inverse kinematic problem [15] that in general, has non-unique solutions and is very difficult to solve because of non-linearity. However, since the NASA robot has a

special configuration, its inverse kinematic problem has a unique closed-form solution.

On the other hand, the forward kinematic problem of the NASA robot does not have a closed-form solution and is solved iteratively. As part of the hardware, six linear voltage differential transformers (LVDT) are mounted parallel to the six linear actuators to measure their extensions. Position feedback is then achieved by taking the actual extension available from measurements and transforming them into the position and orientation of the IEE. This problem is a forward kinematic problem of an open loop robot like the Stanford manipulator.

Fig. 3 shows the general orientation of the manipulator with two coordinate systems, and Fig. 4 illustrates the vector relationships between the origins of each coordinate system and the actuator attachment point of each platform. These relationships yield the vector equations

$$r = A + R \tag{2a}$$

Subtracting (2b) from (2a) and solving for 1 yields

$$\begin{array}{ccc}
1 & = A & + R - B \\
i & i & i
\end{array} \tag{3}$$

where l is defined with respect to the fixed reference frame.

Since the known vectors \mathbf{A} are defined with respect to \mathbf{i} , \mathbf{m}

the moving frame, we must apply an Euler angle transformation in order to get the corresponding vector \mathbf{A}_i in the fixed reference frame. Applying this transformation , we obtain

$$\begin{array}{c}
T \\
1_{i} = [T] A_{i,m} + R - B_{i}
\end{array}$$
(4)

In Equation (4) $A_{i,m}$ and B_i are known constant vectors and R and $[T]^T$ are calculated from the given values of x,y,z,ψ,θ and Φ .

The inverse could be obtained if the vectors $\mathbf{l_i}$ were available, but LVDT measures only the magnitude of the corresponding actuator length $|\mathbf{l_i}|$ and not the required vector. However, because of its closed-loop configuration, the forward kinematics problem of the NASA robot can be solved by solving six nonlinear equations with six unknowns (x,y,z,ψ,θ) and Φ . Newton-Raphson method is employed to solve this problem [16].

The iteration formula has the form

$$\alpha_{n+1} = \alpha_n - \begin{bmatrix} \frac{\partial f(\alpha_n)}{\partial \alpha_n} \\ \frac{\partial \alpha_n}{\partial \alpha_n} \end{bmatrix} = f(\alpha_n)$$
 (5)

Knowing the actual magnitude of the actuators, function $\boldsymbol{f}_{i}\left(\boldsymbol{\alpha}\right)$ can be defined as

$$f_{i}(\alpha) = l_{i}^{T} l_{i} - \left| l_{i} \right|_{a}^{2}$$
 (6a)

where

$$\alpha = \begin{bmatrix} x \\ y \\ z \\ \psi \\ \theta \\ \Phi \end{bmatrix}$$
 (6b)

and

$$f(\alpha) = \begin{bmatrix} f_1(\alpha) \\ f_2(\alpha) \\ \vdots \\ f_6(\alpha) \end{bmatrix}$$
(6c)

and $|1\rangle_{a}$ is the actual magnitude of the actuator i obtained from measurement. Using an appropriate set of initial conditions, Equation (5) is repeated until desired accuracy is met.

The Newton-Raphson's method has been applied in the NASA passive compliant robot system to update the position of the compliant platform relative to the compliant base and to compute the position of the movable platform with respect to the fixed platform [8]. This technique requires ample memory and enormous computation time. Therefore, it is suggested to investigate other Newton-Raphson's method. Some of these methods are known as Quasi-Newton methods [17-18]. The Newton-Raphson's iteration diverges if the initial guess is not close to a correct solution [19]. Using Newton-Raphson-Kantorovich theorem [18] it is possible to find how close the initial

guess should be.Consequently, to assure the convergence of the Newton-Raphson iteration it is necessary to implement the Newton-Raphson-Kantorovich theorem using computer software [10].

5. The 2-DEGREE-OF-FREEDOM ROBOT

In order to examine the performance of proposed hybrid control scheme, we begin our study with the 2-degree-of-freedom robot since the dynamic behaviour of the complete six-degree-of-freedom robot is highly complex.

Fig. 5 illustrates a simplified version of the physical robot consisting of 2 linear actuators, 2 LVDT's, 2 force sensors and an IEE. The actuators are connected to a fixed platform using one-degree-of-freedom gimbals. The LVDT's are mounted parallel to the actuators to measure the lengths l_1 and l_2 . The 2 force sensors are mounted in series with actuators to measure the reaction forces in the direction of the actuator axes.

5.1 KINEMATICS OF THE MANIPULATOR

Manipulator kinematics is a study of the geometry of manipulator arm without regard to the forces which cause it. In this section, the mathematical relations required to describe the system motion and the fundamental equations that govern kinematic behavior are derived.

5.1.1 POSITON ANALYSIS

The position of the point P at the tip of the hand with respect to the X-Y coordinate (Cartesian frame) is known if its horizontal and vertical projections \times , and y are known [10]. However, only the lengths $\begin{pmatrix} 1 & \text{and } l \end{pmatrix}_2$ are known from the LVDT's. From the geometry of Fig. 6, we obtain

and
$$(d-x) + y = l_2^2$$
 (8)

where d is known from the geometry of the system. Since land 1 are variables from measurements, we can rewrite 2 Equations (7) and (8) as

$$f_{1}(x,y) = x + y - 1 = 0$$
 (9)

$$f_{2}(x,y) = (d-x) + y - 1 = 0$$
 (10)

To solve these equations for x and y, we can use iteration techniques such as Newton-Raphson method which was described in section 4. For 2-degree-of-freedom case, the equations can be solved by analitical method. Simplifying Equation (8) as

and substituting (x + y) from Equations (7) into (11), we obtain

$$1_1^2 + d^2 - 2dx = 1_2^2$$
 (12)

Solving for x yields

$$1^{2} - 1^{2} + d^{2}$$

$$= \frac{1}{2} - \frac{2}{2} - \frac{2}{2}$$
(13)

Substituting the value of x into Equation (7)

$$\begin{bmatrix} 1 & -1_{2} & +d & 2 & 2 \\ -1_{2} & -2_{2} & -1_{2} & -1_{2} & 2 & 2 \\ 2d & 2d & 2d & 2d \end{bmatrix}$$
 (14)

and solving for y we get

5.1.2 <u>VELOCITY DERIVATION</u>

A manipulator can, in general, be described as a series of links connected at joints. The angles between the links, called the joint angles, are typical joint variables. However, some type of links (prismatic joints) can grow longer, and in that case the joint variable may actually be the length of the link [6].

The position of all the links of a manipulator of n degrees of freedom can be specified with a set of n joint variables. This set of variables is often referred to as

4 4

the $n\times 1$ joint vector. The space of all such joint vectors is referred to as joint space. There are three representations of a manipulator's position and orientation: descriptions in actuator space, joint space and Cartesian space[15].

In section 5.1.1, we derived the kinematic equations of the hand in Cartesian coordinate system. The forward kinematic transform for this manipulator relates the X-Y position of the hand to the joint variables. For our 2-degree-of-freedom robot with two prismatic actuators, the joint variables can be chosen any of the following sets: $(1_1, \theta_1), (1_2, \theta_2), (\theta_1, \theta_2)$ and $(1_1, 1_2)$. As the LVDT's can measure the lengths $(1_1, 1_2)$, we chose these variables as the joint variables. The inverse kinematic relates the set of joint variables to the desired position and orientation of the hand relative to Cartesian coordinates (Equations (7) and (8)).

In order to move the end-effector in a specified direction at a specified speed, it is necessary to coordinate the motion of the individual joints. For this purpose we drive the differential relationships between the joint displacement and the hand location [3]. The velocity relationship between the joints and hand is determined by the manipulator Jacobian

V=J1. (16a)

Differentiating Equations (13) and (15) with respect to time

COT

and

$$\dot{y} = -\frac{1}{2d} \cdot \frac{1}{2} \cdot \frac{1}{2} \cdot \frac{1}{1} \cdot \frac{1}{1} - 2 \cdot (21_1 \cdot \frac{1}{1} - 21_2 \cdot \frac{1}{2}) \cdot (1_1^2 - 1_2^2 + d^2) \cdot (4d^2 \cdot 1_1^2 + d^2) \cdot (4d^2 \cdot 1_1^$$

$$\left(1_{1}^{2}-1_{2}^{2}+d^{2}\right)^{2}J^{-\frac{1}{2}}$$
 (18)

o r

$$\dot{y} = \frac{1}{4d} \{8d^2 \, \mathbf{1}_1 \, \dot{\mathbf{1}}_1 - 4\mathbf{1}_1 \, (\mathbf{1}_1^2 - \mathbf{1}_2^2 + d^2) \, \dot{\mathbf{1}}_1 + 4\mathbf{1}_2 \, (\mathbf{1}_1^2 - \mathbf{1}_2^2 + d^2) \, \dot{\mathbf{1}}_2 \}$$

$$[4d^2 \, \mathbf{1}_1^2 - (\mathbf{1}_1^2 - \mathbf{1}_2^2 + d^2)^2 \, \mathbf{1}_1^2]$$

$$= \frac{1}{4d} \left[(8d^2 l_1 - 4l_1^3 + 4l_1 l_2^2 - 4l_1 d^2) \right]_1 + 4l_2 (l_1^2 - l_2^2 + d^2) \right]_2$$

$$[4d^{2}l_{1}^{2} - (l_{1}^{2} - l_{2}^{2} + d^{2})^{2}]^{-\frac{1}{2}}$$
(19)

Rewriting the Equations (17) and (19) yield

$$\dot{\hat{x}} = (1_1/d)\dot{1}_1 - (1_2/d)\dot{1}_2$$
 (20)

and

where

$$Q = [4d^21_1^2 - (1_1^2 - 1_2^2 + d^2)^2]^{\frac{1}{2}}$$
 (22)

Writing in vector form

$$\begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix} = \begin{bmatrix} A & A \\ 1 & A \end{bmatrix} \begin{bmatrix} \dot{1} \\ 1 \\ \dot{1} \\ \dot{2} \end{bmatrix}$$
 (23)

Thus the matrix

$$J = \begin{bmatrix} A & A \\ 1 & 2 \\ A & A \end{bmatrix}$$
 (24)

is the manipulator Jacobian,

Where

$$A_1 = 1_1/d$$
, $A_2 = -1_2/d$, $A_3 = [1_1(d^2 + 1_2^2 - 1_1^2)]/dQ$

and

$$A_4 = E1_2 (d^2 + 1_1^2 - 1_2^2) 3/d0$$
.

Note that the elements of the Jacobian are functions of joint displacements, and therefore vary with the arm configuration.

5.1.3 ACCELERATION DERIVATION

The linear acceleration of point P with respect to Cartesian system can be obtained by differentiating Equations (20) and (21) as the following:

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$$\ddot{y} = \frac{2 \cdot 2 \cdot 2}{d\Omega} \cdot \frac{1}{1} \cdot \frac{1}$$

where

$$B_1 = d^2 + l_2^2 - l_1^2$$
 and $B_2 = d^2 + l_1^2 - l_2^2$ (28)

5.2 DYNAMICS OF THE MANIPULATOR

dQ³

The dynamic behavior can be described in terms of the rate of change of the arm configuration in relation to the joint torques exerted by the actuator [3]. Equations of motions, govern the dynamic reponse of the arm linkage to input joint torque. There are two methods to obtain the

equations of motion: Newton-Euler, and Lagrangian formulation.

5.2.1 NEWTON-EULER FORMULATION

The equation of motion is derived by the direct dynamic system interpretation of Newton's Second Law of Motions, which describes dynamic system in terms of force and momentum. The equations include the constraint forces acting between adjacent links which must be eliminated. The dynamic equations of a rigid body can also be represented by two equations:

i)Newton's equation of motion for a mass particle, which is the translational motion of the center of mass (centroid). For a rigid body where center of mass is accelerating with $\overset{\bullet}{V_C}$, the force F, acting at the center of mass is

$$F_{i} = m_{i} \lor_{Ci}$$
 (29)

where, m is the total mass of the body, $V_{\rm Ci}$, the linear velocity of the centroid of link i with respect to the base coordinates (inertial reference frame), and $m_{\rm i} \dot{V}_{\rm Ci}$, inertial force [15].

ii) Euler's equation of motion, which is the rotational

motion of the centroid. For a rigid body rotating with angular velocity \mathbf{w}_1 and angular acceleration $\dot{\mathbf{w}}_2$, the moment \mathbf{N}_1 , which must be acting on the body to cause this motion is

$$N=I_{c}\dot{w} + w \times I_{c}w \tag{30}$$

where I is the inertia moment of the body, I $\overset{\bullet}{\text{c}}$ inertia torque, and I w is the angular momentum.

5.2.2 LAGRANGIAN FORMULATION

In the Lagrangian method, the system dynamic behavior is described in terms of work and energy, using generalized coordinates. All the workless forces and constraint forces are automatically eliminated in this method. The resultant equations are generally compact and provide a closed-form expression in terms of joint torques and joint displacements. Furthermore, the derivation is simpler and more systematic than in the Newton-Euler approach. As in kinematics, we need to solve the inverse problem of finding the necessary input torques to obtain a desired output motion (inverse dynamics).

Let q,, q be generalized coordinates, T and P

the total kinetic and potential energy stored in the

dynamic system. We define the Lagrangian "L" by

$$L(q, \dot{q}) = T - P \tag{31}$$

CO"

$$d/dt(aL/aq_i) - aL/aq_i = 0 \qquad i=1,...n \qquad (32)$$

where \mathbf{Q}_{i} is the generalized forces corresponding to the generalized coordinates \mathbf{q}_{i} .

5.2.3 EQUATIONS OF MOTION

To formulate the equations of motion of the robot, we employ the Lagrangian and Newton-Euler approaches.

i) Lagrangian method

Fig. 6 represents the forces acting on the robot system. The robot task is to keep in contact with a reaction surface while maintaining a constant applied force \mathbf{f}_0 . This task occurs for example in the process of painting a surface. The robot is moving in a horizontal plane (to the right), while maintaining a constant vertical force on the reaction table. For simplicity, we assume that the table is smooth enough so that the table friction is negligible, joint and actuator friction are also negligible. Therefore, the reaction force will be normal to the hand ($\mathbf{f}_{\mathbf{v}}$).

Using the Lagrangian formulation, we will derive the equations in terms of the two independent joint variables,

the length l_{ij} and the angle θ_{ij} [21].

We begin by computing the velocity of the centroids $^{\text{C}}$, $^{\text{C}}$ and the end-effetor F. From geometry, the Cartesian positions are

$$P = 1 \quad \text{cos} \quad \theta_1 \tag{33}$$

$$y_{\mathbf{p}} = -1_{\mathbf{1}} \sin \theta_{\mathbf{1}} \tag{34}$$

$$\underset{\text{Cl}}{\overset{=(1}{1}}/2)\cos\theta$$
 (35a)

and
$$y_{cl} = -(1_1/2)\sin\theta_l$$
 (35b)

Similarly,

and
$$y_{C2} = -(\frac{1}{2})\sin\theta_2$$
 (36b)

To eliminate θ_2 and θ_2 from the equations, we can write

$$_{1}^{1} \cos \theta_{1}^{=d-1} \cos \theta_{2} \tag{37}$$

and
$$-1_1 \sin \theta_1 = -1_2 \sin \theta_2$$
 (38)

Substituting from Eqs. (37) and (38) into (36a) and (36b) we get

and
$$y_{C2} = -1_1 \sin \theta_1 / 2 \tag{40}$$

Differentiating the above equations with respect to time yields the Cartesian velocities as

$$V_{\mathbf{p}} = \begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix} = \begin{bmatrix} \mathbf{1}_{1} \cos \theta_{1} - \mathbf{1}_{1} \sin \theta_{1} \dot{\theta}_{1} \\ -\mathbf{1}_{1} \sin \theta_{1} - \mathbf{1}_{1} \cos \theta_{1} \dot{\theta}_{1} \end{bmatrix} = \begin{bmatrix} \cos \theta_{1} & -\mathbf{1}_{1} \sin \theta_{1} \\ & & \\ \sin \theta_{1} & -\mathbf{1}_{1} \cos \theta_{1} \end{bmatrix} \begin{bmatrix} \mathbf{1}_{1} \\ \dot{\theta}_{1} \end{bmatrix}$$
(41)

and

$$V_{C1} = V_{C2} = \begin{bmatrix} (1_{1} \cos \theta_{1}^{-1} \sin \theta_{1} \theta_{1}^{*})/2 \\ (1_{1} \sin \theta_{1}^{+1} \cos \theta_{1} \theta_{1}^{*})/2 \end{bmatrix}$$

$$(42)$$

The magnitude of the velocity vectors can now be computed as

$$V_{C1}^2 = V_{C2}^2 = (x_1)^2 + (y_1)^2$$
 (43a)

CO.

$$V_{C1}^{2} = V_{C2}^{2} = 1/4E(\hat{i}_{1}\cos\theta_{1})^{2} + (\hat{i}_{1}\sin\theta_{1}\hat{\theta}_{1})^{2} - 2\hat{i}_{1}\hat{i}_{1}\sin\theta_{1}\cos\theta_{1}\hat{\theta}_{1}) + 1/4E(\hat{i}_{1}\sin\theta_{1}^{2} + (\hat{i}_{1}\cos\theta_{1}\hat{\theta}_{1})^{2} + (\hat{i}_{1}\cos\theta_{1}\hat{\theta}_{1})^{2} + 2\hat{i}_{1}\hat{i}_{1}\sin\theta_{1}\cos\theta_{1}\hat{\theta}_{1})$$
(43b)

$$v_{C1}^{2} = v_{C2}^{2} = (1_{1} + 1_{1} + 1_{1} + 1_{1}) / 4 . \tag{43c}$$

Similarly,

The kinetic energy of a mass m with linear velocity V and angular velocity w is

$$T=1/2 \text{ mV}^2 + 1/2 \text{ Iw}^2$$
 (45a)

Note that as we assumed point mass, then the moment of inertia ${\rm I}_1{=}{\rm I}_2{=}0$, thus ${\rm Iw}^2{=}0$.

Therefore.

$$T=\Sigma(1/2 \text{ m}_{i}^{V})$$
 $i=1,...,N$ (46a)

$$T=1/2 \text{ M V}^{2} + 1/2 \text{ M V}^{2} + 1/2 \text{ M V}^{2} + 1/2 \text{ M V}^{2}$$

$$(46b)$$

where M , M and M are the masses of the two linear actuators and the IEE, respectively. Substituting the values of the velocities into (46b) yields

$$T = \left[\left(\frac{M_1 + M_2}{1} \right) / 8 + \frac{2}{3} \right] \left(\frac{2}{1} + \frac{2}{1} + \frac{2}{1} \right)$$

$$(46c)$$

The actuators are assumed to have identical masses, $M_1 = M_2 = M$. This assumption changes (46c) to

$$T = (M/4 + M_3/2) \begin{pmatrix} 1 & 2 & 2 & 2 \\ 1 & +1 & \theta & 1 \end{pmatrix}.$$
 (46d)

Now , we find the potential energy, using the formula

$$F=mgh$$
 (47a)

where h is the height and g the acceleration of gravity. Thus,

$$F_1 = Mg(-1\sin\theta_1/2) \tag{47b}$$

$$P_2 = Mg(-l_2 \sin\theta_2/2) \tag{48}$$

$$P_{3} = M_{3}g(-l_{1}\sin\theta_{1}) \tag{49}$$

where from Eq. (38)

$$l_1 \sin \theta_1 / 2 = l_2 \sin \theta_2 / 2 \qquad (49a)$$

The total potential energy is

$$F = F_1 + F_2 + F_3 = -(M + M_3) g l_1 \sin \theta_1$$
 (50)

We now combine Equations (46d) and (50) to produce the Lagrangian for this manipulator

The free-body diagram of the Link 1 is shown in Fig. 7.

The total force applied on the system is

$$F = f_{m1} - f_{m2} \cos \theta - f_{y} \sin \theta_{1}$$
 (52)

where
$$\theta = \pi - (\theta_1 + \theta_2)$$
. (53)

Therefore,

$$F = f_{m1} + f_{m2} \cos(\theta_1 + \theta_2) - f_y \sin\theta_1$$
 (54)

Although there is not any torque about θ_1 , but when the second link is extending or contracting, it makes the first one to rotate. So we can state that the forces f_{m2} and f_y apply a torque about θ_1 as the following:

$$T_{\theta} = f_{m2} l_{1} \sin \theta + f_{y} l_{1} \cos \theta_{1}$$
 (55)

Cir

$$T_{\theta} = [f_{m2} \sin(\theta_1 + \theta_2) + f_y \cos(\theta_1)]$$
 (56)

In Equations (55) and (56) f is pushing and f is ml m2 pulling. Using Equation (32) we can find the forces applied by the linear actuators by considering

$$\frac{\partial L}{\partial \hat{1}} = (M/2 + M_3) \hat{1}_1$$
 (57)

$$\frac{d}{dt} = \frac{\partial L}{\partial i} = (M/2 + M_3) \frac{i}{1}$$
(58)

$$\frac{\partial L}{\partial l_1} = (M/2 + M_3) l_1 \cdot \frac{e_1^2}{e_1^2} + (M + M_3) g \sin e_1.$$
 (59)

Therefore,

$$F = (M/2+M_3) \frac{1}{1} - (M/2+M_3) \frac{1}{\theta_1} - (M+M_3) g \sin \theta_1$$
 (60)

Similarly, from the torque

$$\frac{\partial L}{\partial \theta_{1}} = (M/2 + M_{3}) 1 \frac{\partial \theta_{1}}{\partial \theta_{1}} \tag{61}$$

$$\frac{d}{dt} = \frac{\partial L}{\partial \theta_{1}} = (M/2 + M_{3})(2l_{1} l_{1} \theta_{1} + l_{1} \theta_{1})$$

$$(62)$$

$$3L$$
---- = $(M+M)$ gl $cos \theta$. (63)
 θ_1

Thus,

$$T_{\theta} = (M/2 + M_3)(21_1 \cdot 1_1 \cdot \theta_1 + 1_1 \cdot \theta_1) - (M + M_3)g1_1 \cos \theta_1$$
 (64)

Substituting the values of F and T from (54) and (56) into Equations (60) and (64) yield

$$f_{m1} + f_{m2} \cos(\theta_1 + \theta_2) - f_y \sin \theta_1 = (M/2 + M_3) \frac{1}{1} - (M/2 + M_3) \frac{2}{1} \theta_1 - (M/2 + M_3) \frac{2}{1} \theta_1$$

$$(M + M_3) g \sin \theta_1$$
(65)

$$f_{m2} = in(\theta_1 + \theta_2) + f_y = cos\theta_1 = (M/2 + M_3) l_1\theta_1 + (M/2 + M_3) (2\dot{1}\dot{\theta}_1) - (66)$$

To write the above equations in terms of joint variables ${\bf l}_1$ and ${\bf \theta}_1$, we must eliminiate ${\bf \theta}_2$. From trigonometry we obtain

$$\cos(\theta_1 + \theta_2) = \cos\theta_1 \cos\theta_2 - \sin\theta_1 \sin\theta_2 \tag{67}$$

Substituting the values of $\cos\theta_2$ and $\sin\theta_2$ from Equations (37) and (38) into (67) we get

$$\cos(\theta_1 + \theta_2) = \cos\theta_1 [(d-1_1 \cos\theta_1)/1_2] - \sin\theta_1 [(l_1 \sin\theta_1)/l_2]$$
 (68)

or
$$\cos(\theta_1 + \theta_2) = (\cos\theta_1 - l_1)/l_2$$
. (69)

Similarly,

$$\sin(\theta_1 + \theta_2) = \sin\theta_1 \cos\theta_2 + \cos\theta_1 \sin\theta_2 \tag{70}$$

C+3"

$$\sin(\theta_1 + \theta_2) = \sin\theta_1 \left[(d-l_1 \cos\theta_1)/l_2 \right] + \cos\theta_1 \left[(l_1 \sin\theta_1)/l_2 \right]$$

$$= d \sin\theta_1/l_2$$
(71)

From Fig. 6, we obtain

$$l_2 = (l_1^2 + d^2 - 2dl_1 \cos \theta_1)^{\frac{1}{2}}$$
 (73)

Substituting l_2 into Equations (69) and (71)

and
$$\sin(\theta_1 + \theta_2) = ----$$
 (75)
$$2 \quad 2 \quad \frac{1}{2} + d \quad -2d1 \cos \theta_1$$

Equations (65) and (66) can be rewritten as

$$f_{m1} + f_{m2} = (t) - f_{sin\theta} = k_1 l_1 - k_1 l_{\theta_1} - k_2 gsin_{\theta_1}$$
 (76)

$$f_{m2} b(t) + f_{cos\theta_1} = K_1 \theta_1 + 2K_1 i_1 \theta_1 - K_2 gcos\theta_1$$
(77)

where

$$K_1 = M/2 + M_3$$
 ; $K_2 = M + M_3$ (78)

 $a(t) = (d\cos\theta_1 - l_1)/c(t) + b(t) = d\sin\theta_1/c(t);$

and f_{ml} and f_{m2} are actuator forces. The force f_y is the reaction force produced through contact with the reaction surface and is modelled as follows:

$$f_{y} = \begin{cases} -k_{E}(y-y_{E}) & \text{when in contact} & (y-y_{E}) < 0 \\ \\ 0 & \text{otherwise} & (y-y_{E}) \ge 0 \end{cases}$$

where y $_{\rm E}$ and K $_{\rm E}$ denote the position and stiffness of the reaction surface, respectively [3]. The consant K $_{\rm E}$ is

typically of the order of 10 N/m for a stiff environment.

The term $2K_1 l_1 \theta_1$ in Eq. (77) represents the Coriolis force acting on Link 2 due to the Coriolis effect (when a mass particle m moves at a velocity of V relative to a moving coordinate frame rotating at an angular velocity w, the mass particle has the so-called Coriolis force given by 2m(wxV)). The last term $K_2 g \sin \theta_1$ and $K_2 g \cos \theta_1$ in Equations (76) and (77) accounts for the effect of gravity. The second term on the right side of Eq. (76) which is proportional to the square of the joint velocities shows the centrifugal force.

ii) Newton-Euler Method

In order to reduce the computation time for the simulation of the robot system, the equations of motion in terms of the Cartesian variables X and Y are needed. Fig. 8 illustrates the free body diagram of each link of the robot system. Applying Newton formulation for IEE, we obtain

$$F_{1x} + F_{2x} = M_3 \ddot{x}$$
 (81)

$$-F_{1y} + F_{2y} - M_{3g} + F_{y} = M_{3y}$$
 (82)

Knowing that $\dot{V}_{C1} = \dot{V}_{C2} = \dot{V}_{P}/2$, for Links 1 and 2 we can write

$$-f \sin_{\theta} +F -Mg = My/2$$
ml ly (84)

$$f_{m2} = -F_{m2} = \frac{n}{2x}$$
(85)

$$f_{ml} (y/l) - f_{m2} (y/l_2) + f_{=K} y + K_3 g$$
 (88)

where $K_3 = 2M + M_3$ (89)

and $\begin{pmatrix} 1 & 1 & 1 \\ 1 & 2 \end{pmatrix}$ and y are related through Equations (7) and (8).

5.3 HYBRID CONTROL OF THE 2-DEGREE-OF-FREEDOM ROBOT

The block diagram of the hybrid control system is shown in Fig. 9. Both position and velocity feedback are employed to improve the response. We assume that the robot is to move from the initial position $x_i = A$, $y_i = B$ to the final position $x_i = A$, while exerting a force $f_i = f_i$ in

the vertical direction. Manipulator is moving in a horizontal plane so that gravity may be ignored [6]. The desired velocity in X direction will be $\stackrel{\circ}{x}$ =0 while in y direction is zero.

5.3.1 POSITION AND VELOCITY CONTROL

For a 2-degree-of-freedom robot the selection matrix defined in Eq.(1), will be

$$Sjj = \begin{bmatrix} 1 & O \\ O & 1 \end{bmatrix}$$

Since we are interested in controlling position in x direction (s =0), and force in Y direction (s =1), the corresponding compliance matrix is $S = \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix}$

and
$$[I]-[S] = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} - \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}$$
 (90)

where [I] is the identity matrix. The position error is

$$\Delta X(t) = X_{\mathbf{d}} - \Lambda(\mathbf{q}(t)) = \begin{bmatrix} x_{\mathbf{d}} \\ y_{\mathbf{d}} \end{bmatrix} - \begin{bmatrix} x(t) \\ y(t) \end{bmatrix} = \begin{bmatrix} A \\ B \end{bmatrix} - \begin{bmatrix} x(t) \\ y(t) \end{bmatrix} = \begin{bmatrix} A - x(t) \\ B - y(t) \end{bmatrix}$$
(91)

This error is mapped into the position controlled subspace. The compliance matrix S serves as a mask to separate the force controlled and position controlled degrees of freedom. The Cartesian error signals are obtained as

$$X_{\mathbf{e}}(\mathbf{t}) = \begin{bmatrix} x_{\mathbf{e}}(\mathbf{t}) \\ y_{\mathbf{e}}(\mathbf{t}) \end{bmatrix} = \{ [] - [S] \} \Delta X(\mathbf{t}) = \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} A - x(\mathbf{t}) \\ B - y(\mathbf{t}) \end{bmatrix} = \begin{bmatrix} A - x(\mathbf{t}) \\ 0 \end{bmatrix}$$
(92)

Transforming them back into the joint space, we find the corresponding joint variables errors. Since X_e is a differential change, we can use the Jacobian to find the corresponding (approximate) differential change in the vector of joint variables

$$q(t) = \begin{bmatrix} 1_1(t) \\ 1_2(t) \end{bmatrix}$$

Recalling J from Equation (24) we find its inverse as follows:

where
$$\begin{bmatrix} 2 & 2 & 2 \\ 1_2 (d + 1_1 & -1_2) & 1_2 \\ ------ & --- \\ dQ & d \end{bmatrix}$$
Adj $|J| = \begin{bmatrix} 2 & 2 & 2 \\ 1_1 (d + 1_2 - 1_1) & 1_1 \\ ----- & --- \\ dQ & d \end{bmatrix}$ (93b)

CONT

$$Det |J| = \frac{2l_1}{0} \frac{l_2}{0}$$
 (93e)

Therefore,

For simplicity, we define

$$J^{-1} = \begin{bmatrix} E_1 & E_2 \\ E_3 & E_4 \end{bmatrix}$$
 (95)

where E_i 's are given in Equation (94). Thus

$$q_{e}(t) = \begin{bmatrix} l_{e1}(t) \\ l_{e2}(t) \end{bmatrix} = J^{1} X_{e}(t)$$
 (96)

$$q_{e}(t) = \begin{bmatrix} E_{1} & E_{2} \\ E_{3} & E_{4} \end{bmatrix} \begin{bmatrix} A + x(t) \\ O \end{bmatrix} = \begin{bmatrix} E_{1}[A - x(t)] \\ E_{3}[A - x(t)] \end{bmatrix}$$
(97a)

The error in joint variables will be

$$\begin{array}{ccc}
1 & (t) = E_1 [A - x(t)] \\
e1
\end{array} \tag{97b}$$

$$\begin{array}{ccc}
1 & (t) = E_3[A - x(t)] \\
e2
\end{array} \tag{97c}$$

Applying the same method for velocity control, we find

$$\Delta \dot{\mathbf{x}}(t) = \dot{\mathbf{x}}_{\mathbf{d}} - \Lambda(\dot{\mathbf{q}}(t)) = \begin{bmatrix} \dot{\mathbf{x}}_{\mathbf{d}} \\ \dot{\mathbf{y}}_{\mathbf{d}} \end{bmatrix} - \begin{bmatrix} \dot{\mathbf{x}}(t) \\ \dot{\mathbf{y}}(t) \end{bmatrix} = \begin{bmatrix} 0 & -\dot{\mathbf{x}}(t) \\ 1 & \\ 0 & -\dot{\mathbf{y}}(t) \end{bmatrix}$$
(98)

$$\dot{\mathbf{x}} (t) = \begin{bmatrix} \dot{\mathbf{x}} (t) \\ \mathbf{e} \\ \dot{\mathbf{y}} (t) \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} \mathbf{0} - \dot{\mathbf{x}}(t) \\ 1 \\ 0 - \dot{\mathbf{y}}(t) \end{bmatrix} = \begin{bmatrix} \mathbf{0} - \dot{\mathbf{x}}(t) \\ 1 \\ 0 \end{bmatrix}$$
(99)

Therefore, the error in joint velocities will be

To move the joints and reduce the error, we use (proportional derivative) controller for position.

$$u = q(t)[K]$$
 (101a)

$$v = q (t)[K]$$
 (101b)

The error in lengths due to the error in position and velocity are

$$u_{p} = u_{p} = u_{p}$$
 (102a)

For Link 1 we get

and for Link 2

$$u (t)=1 (t)K +1 (t)K (102c)$$

 $rp2 e2 rp2 e2 rpd2$

Substituting the values of l , l , l , and l el e2 el e2 into (102b) and (102c), we obtain

$$u_{rp2}$$
 (t)=E (A-x(t))K +E (Q -x(t))K (103b)

5.3.2 FORCE CONTROL

Knowing the desired force

$$F_{d} = \begin{bmatrix} f_{dx} \\ f_{dy} \end{bmatrix} = \begin{bmatrix} 0 \\ f_{0} \end{bmatrix}$$

and the actual force $F(t) = \begin{bmatrix} f_x \\ f_y \end{bmatrix}$, the force error is computed as

$$\Delta F(t) = F_{d}(t) - F(t) = \begin{bmatrix} 0 \\ f_{0} \end{bmatrix} - \begin{bmatrix} f_{x}(t) \\ f_{y}(t) \end{bmatrix} = \begin{bmatrix} -f_{x}(t) \\ f_{0} - f_{y}(t) \end{bmatrix}$$
(104)

Note that the coordinate system is parallel to the constraint frame. So, the force transformation matrix is

$$\begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

The Cartesian error signal is computed by

$$F_{e}(t) = \begin{bmatrix} f_{xe} \\ f_{ve} \end{bmatrix} = [S]\Delta F(t)$$
 (105a)

$$F_{e}(t) = \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} -f_{x}(t) \\ f_{0} - f_{y}(t) \end{bmatrix} = \begin{bmatrix} 0 \\ f_{0} - f_{y}(t) \end{bmatrix}$$
(105b)

The corresponding Joint space error is found by

$$\tau_{e}(t) = J^{T} F_{e}(t) \tag{106}$$

where $\tau_e(t) = \begin{bmatrix} f_{rel} \\ f_{re2} \end{bmatrix}$ and J is the Jacobian transpose that is

derived as

For simplicity, we define
$$J = \begin{bmatrix} G_1 & G_2 \\ G_3 & G_4 \end{bmatrix}$$
 (108)

then

$$\tau_{e}(t) = \begin{bmatrix} G_{1} & G_{2} \\ G_{3} & G_{4} \end{bmatrix} \begin{bmatrix} O \\ f_{0} - f_{y} \end{bmatrix} = \begin{bmatrix} G_{2}(f_{0} - f_{y}) \\ G_{4}(f_{0} - f_{y}) \end{bmatrix}$$
(109a)

COT

$$f_{rel}(t) = G_2(f_0 - f_y)$$
 (109b)

$$f_{re2}(t) = G_4 (f_0 - f_y)$$
 (109c)

For force control, we will use PE (proportional error) controller. Therefore, the error in lengths due to the error in force will be

$$u_{rf} = K_{rf}^{f}_{re}$$
 (110a)

COT

$$u_{rfl}(t) = K_{rfl} f_{rel}(t) = K_{rfl} G_2(f_0 - f_y(t))$$
 (110b)

$$u_{rf2}(t) = K_{rf2} f_{re2} (t) = K_{rf2} G_4 (f_0 - f_y(t))$$
 (110c)

Finally the inputs to the linear actuators are computed by

$$G_4(f - f (t)) \times (112)$$

0 y rf2

5.3.3 LINEAR ACTUATOR DYNAMICS

Since no transient analysis of DC linear motor is available, we assume that its dynamics is identical to that of DC motor. By neglecting the armature inductance and the load torque, the transfer function between the motor actuating force and the input voltage e_a is obtained by [27]

where we assumed that F(S)=K(T(S)) and

R =armature resistance

J =rotor inertia of motor

K =back emf constant

K =torque constant

B =viscous frictional coefficient

6. SIMULATION

In order to examine the behavior of the proposed hybrid controller, we conducted several simulation runs of the system. The robot dynamics and the hybrid control scheme are simulated using the System Simulation Language (SYSL) developed by E Consulting for the IEM FC.

6.1 RESULTS AND DISCUSSION

Fig. 10 shows the responses of the system that was controlled to stay at a fixed position X =10 cm, Y =-10 cm. while mainting a constant force F =90 N. The responses had small overshoots in vertical force and position, which is caused by the initial forces of the motors (f_{ml} and f_{m2} =6N). The settling time was about 0.1 sec. The error in the horizontal position was about 10 cm and was not recorded in the graph.

Figures 11 and 12 depict the responses of the robot controlled to move from X=0 to X=1 cm horizontally while keeping a constant force F=90 N on the reaction surface. The response shows that the robot reached the destination in 0.5 sec in an "overdamped" fashion. The average velocity in Fig. 11 is 8 cm/sec and in Fig. 12 7cm/sec. The first case shows higher overshoot but faster

response. The vertical force and position were settled after about 0.1 sec with relatively small overshoots.

Figures 13 and 14 demonstrate the control of force in the presence of position disturbances. The controller attempts to maintain a constant 90 N force at a fixed horizontal position on the reaction table while the table is moving away from the robot in y direction at a constant rate 0.1 cm/sec (Fig. 13), and 0.08 cm/sec (Fig.14). The quality of the force was quite good while there was some disturbance in the horizontal position (about 0.01 cm error). When the surface stops moving (after 1 sec), control returns to a stable steady state. At the end of the ramp where there are acceleration there was a little disturbance in the force. In these figures, we note that the vertical position follows the movement of the table closely.

In all simulation runs, the controller gains, actuator parameters and robot masses are adjusted until a stable response with little overshoots is achieved. Controller feedback gains and the numerical values of parameters used in simulation are given below

$$K = 39.$$
 N/m $K = 25.$ N/m $K = 25.$ N/m $K = 10.3$ $K = 10.3$

$$K = 9 \times 10^{5}$$
 N/m $d=23$ cm $1 = 1 = 15$ cm

7. CONCLUSIONS

In this report the hybrid control was proposed to control the compliance of the existing NASA robot. A study of coordinate transformation of the IEE showed that its forward kinematic problem does not have a closed-form solution and consists of non-linear equations which Newton-Raphson's method was employed to solve it.

The kinematics, dynamics and hybrid controller of the 2-degree-of-freedom robot were presented and its simulation results were discussed.

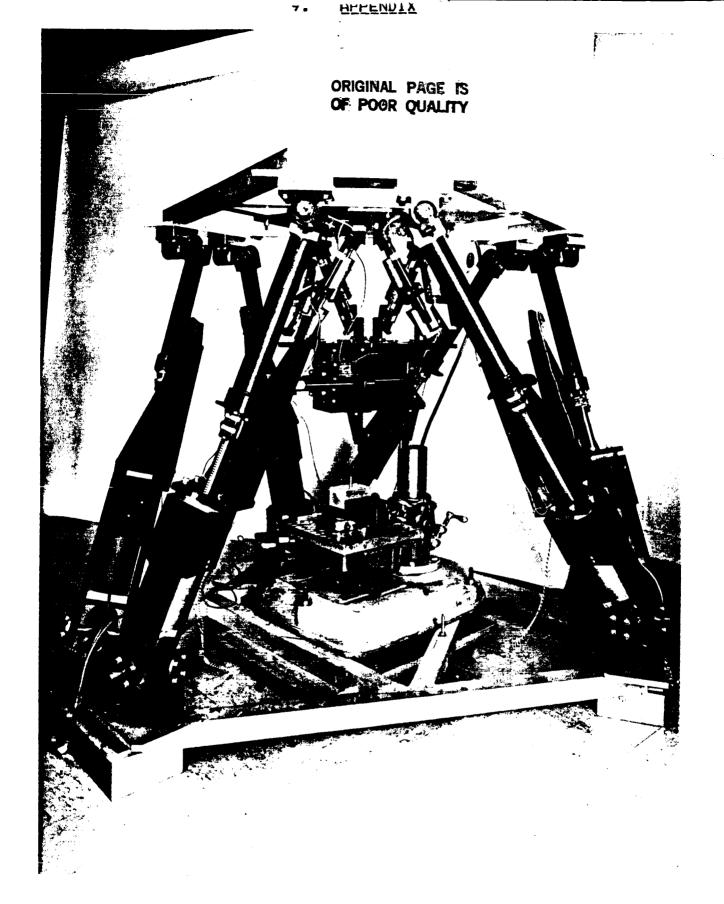
Even though we obtained some good results, no clear relationship between the controller gains and system responses have been obtained. Further research can focus on the development of a guideline to select controller gains for stable and responsive robot systems.

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 $\underline{\text{Fig. 1A}}$: The NASA Passive Compliant Robot System

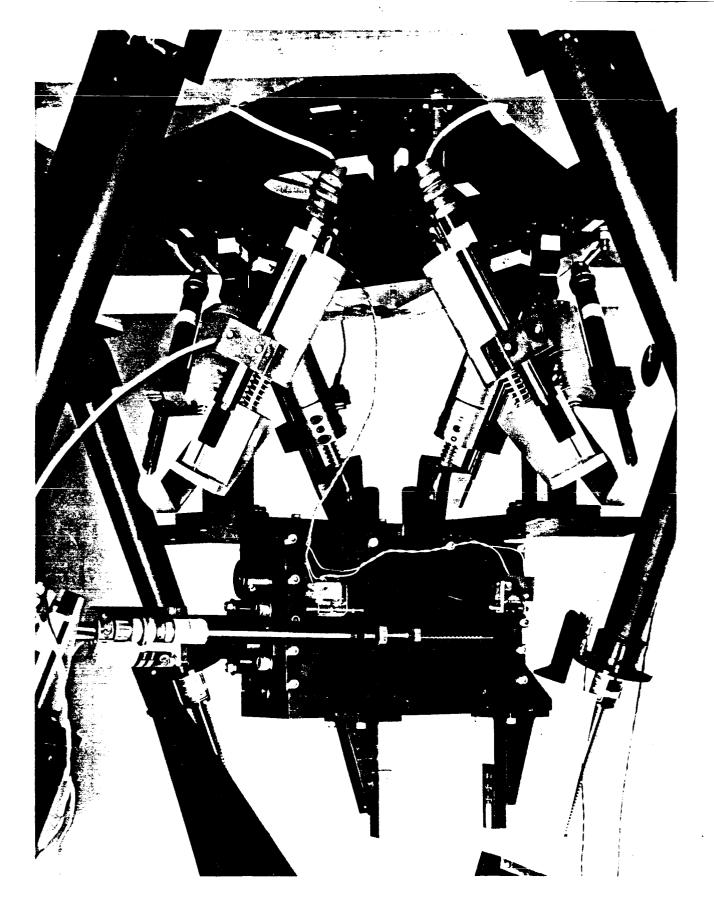


Fig. 1B : The Passive Compliant Platform

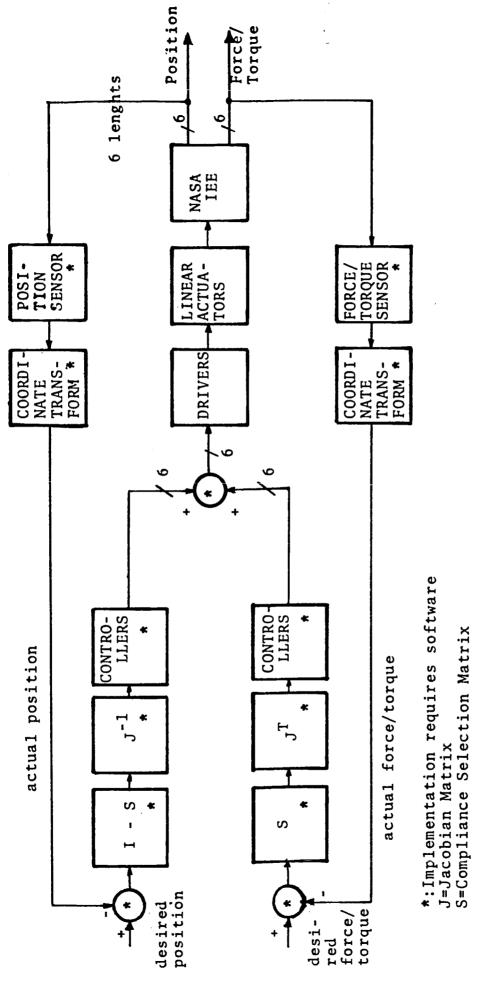


Fig. 2: The Hybrid Control Scheme

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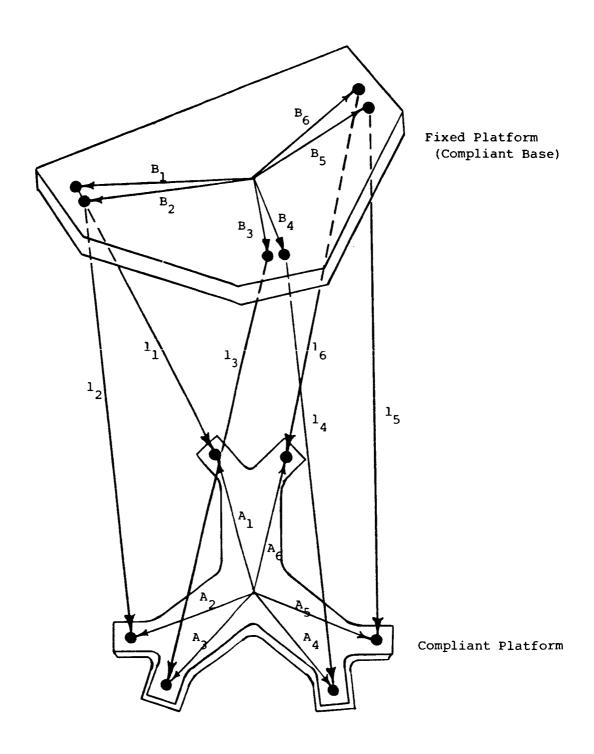


Fig. 3: Particular Orientation of Robot

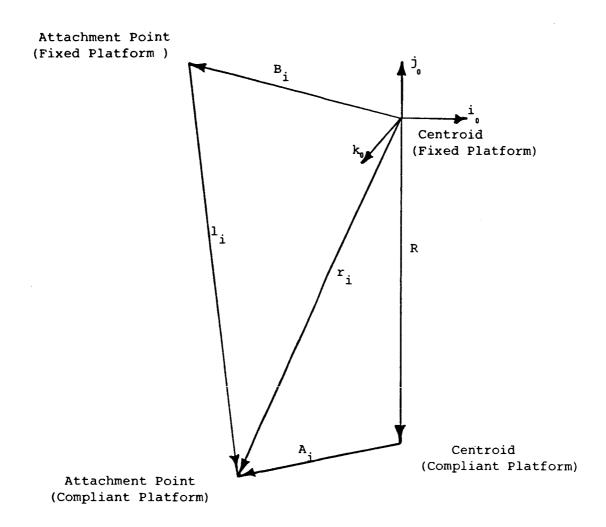


Fig. 4: Vector Relationships for Actuator i

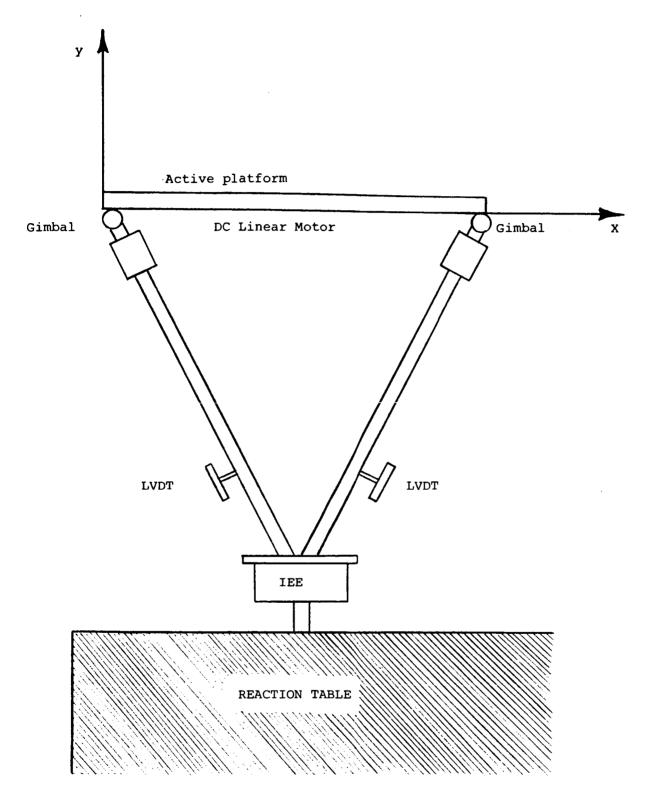


Fig. 5: The 2-Degree-of-Freedom Robot

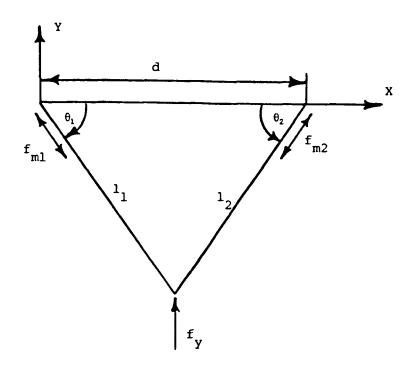
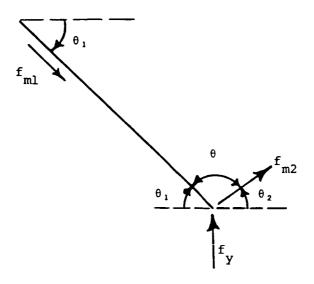


Fig. 6: Forces Acting on the Robot



 $\underline{\text{Fig. 7}}$: The Free-Body Diagram of Link 1

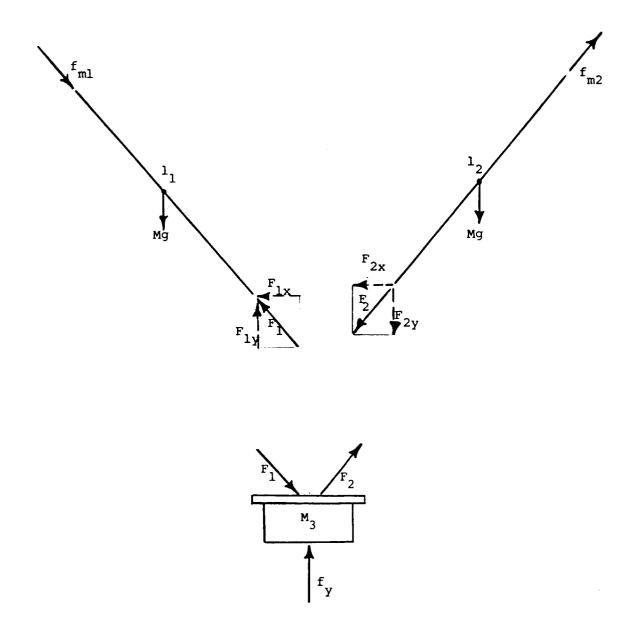


Fig. 8: The Free-Body Diagram

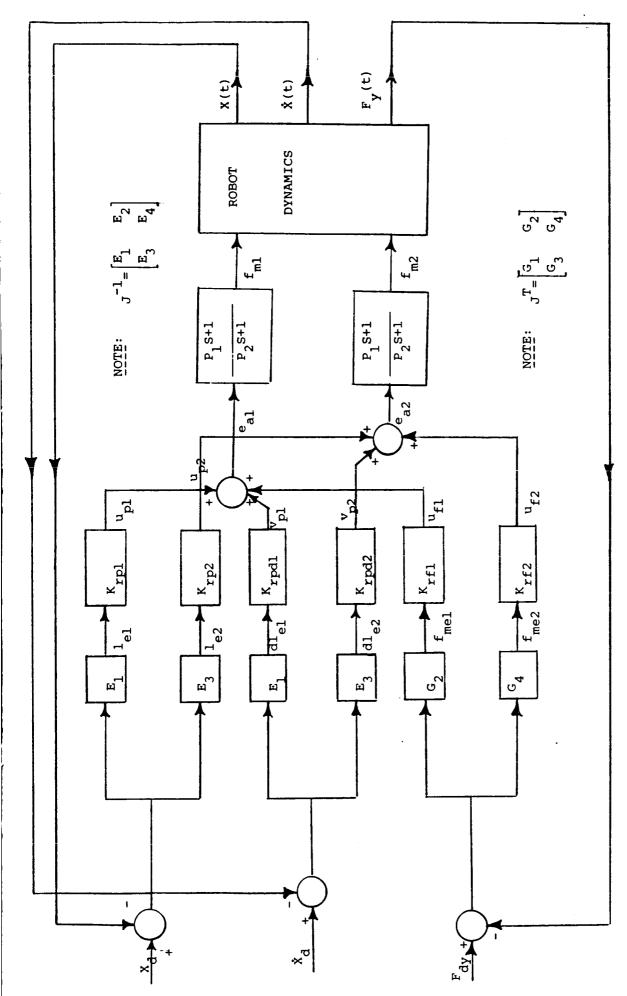


Fig. 9: Implementation of the Hybrid Control Scheme for the 2-Degree-of-Freedom Robot.

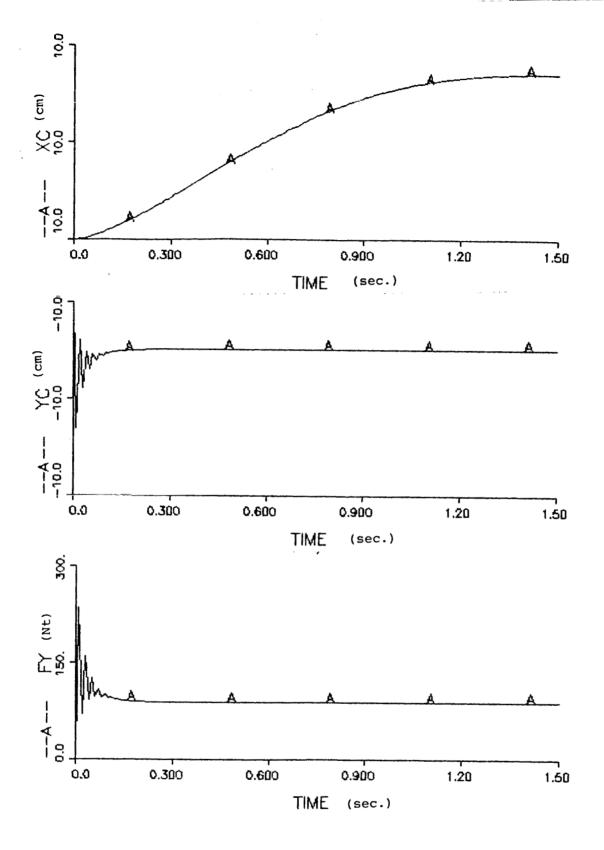


Fig. 10: Responses of the Robot for Control of Constant Force and Fixed Position

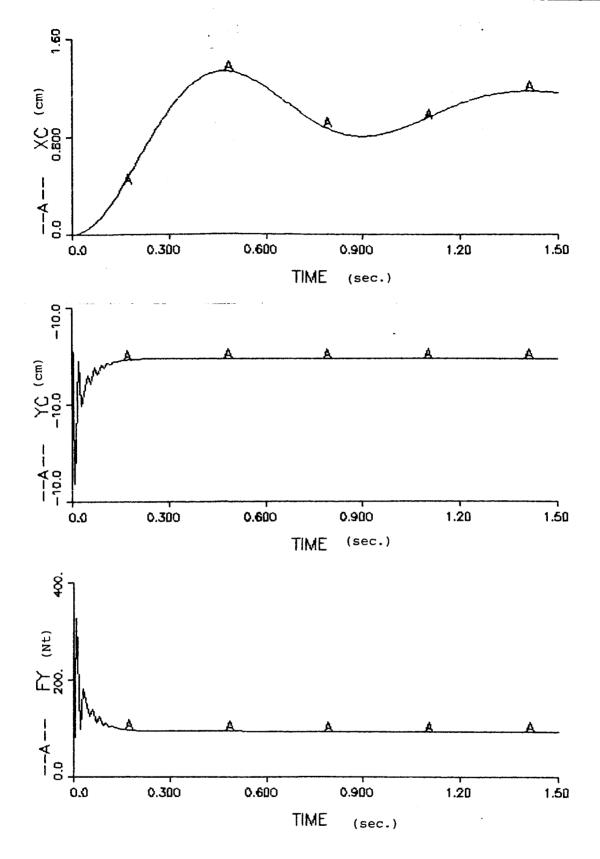


Fig. 11: Responses for Control of Constant Force while Moving on Horizontal Plane (V= 8. cm/sec)

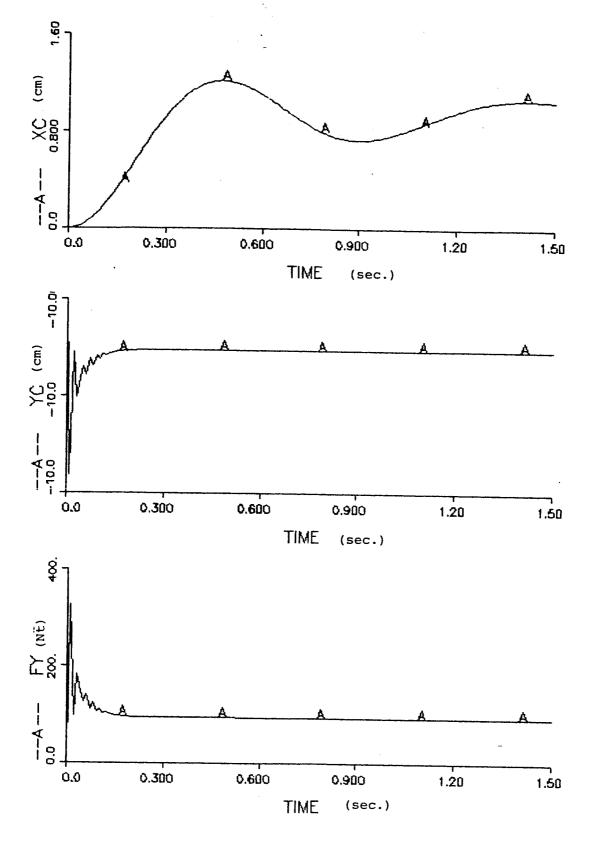


Fig. 12: Responses for Control of Constant Force while Moving on Horizontal Plane (V= 7.cm/sec)

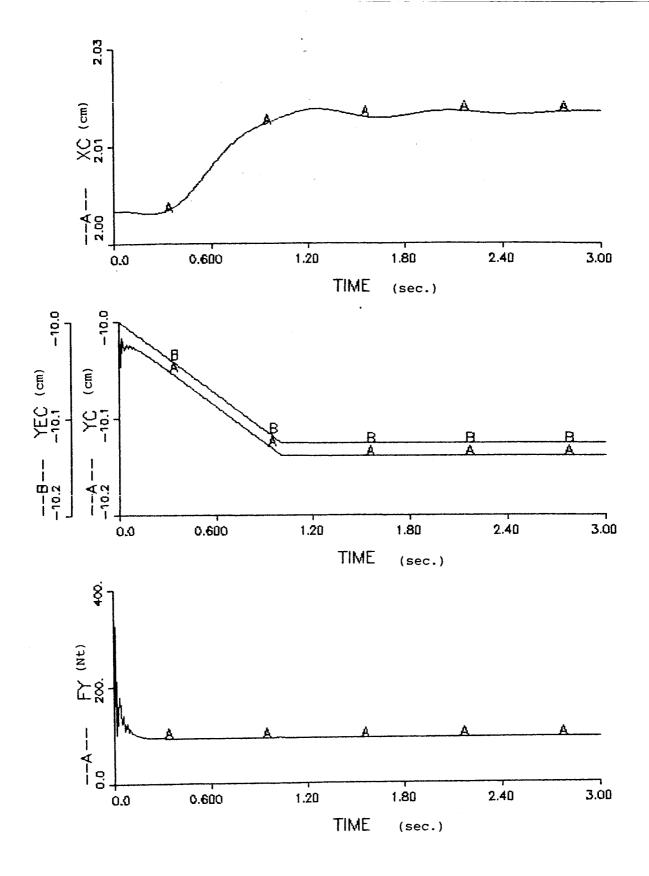


Fig. 13: Responses for Control of Constant Force on a Moving Table ($V_{table} = 0.1 \text{ cm/sec}$)

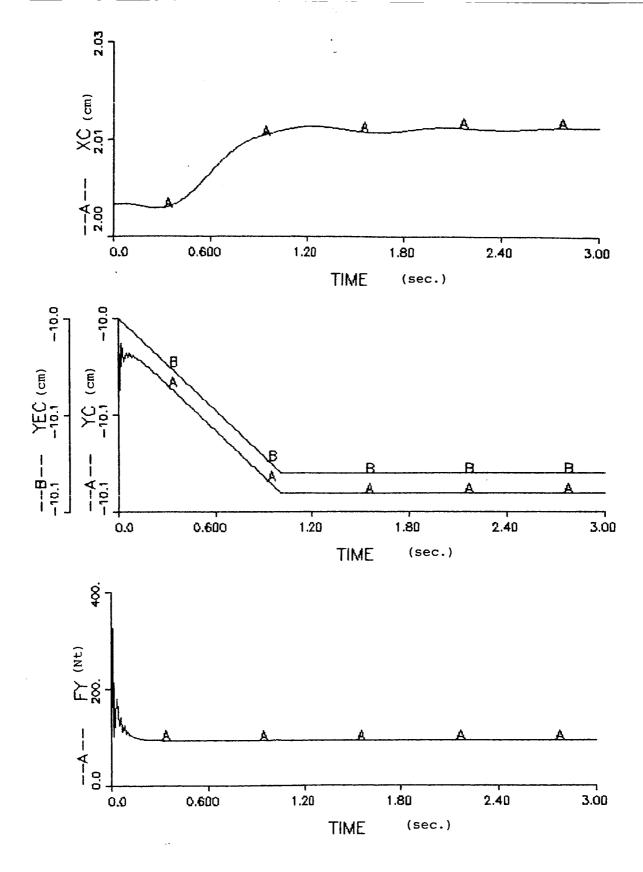


Fig. 14: Responses for Control of Constant Force on a Moving Table ($V_{table} = 0.08 \text{ cm/sec}$)